DOI: 10.37863/umzh.v74i12.6734

UDC 517.9

M. Kiss¹ (Inst. Math. Budapest Univ. Technology and Economics, Hungary)

AN AMBARZUMIAN TYPE THEOREM ON GRAPHS WITH ODD CYCLES ² ТЕОРЕМА ТИПУ АМБАРЦУМЯНА ДЛЯ ГРАФІВ З НЕПАРНИМИ ЦИКЛАМИ

We consider an inverse problem for Schrödinger operators on a connected equilateral graph G with standard matching conditions. The graph G consists of at least two odd cycles glued together at a common vertex. We prove an Ambarzumian-type result, i.e., if a specific part of the spectrum is the same as in the case of zero potential, then the potential must be equal to zero.

Розглянуто обернену задачу для операторів Шредінгера на зв'язному рівносторонньому графі G зі стандартними умовами узгодження. Граф G складається принаймні з двох непарних циклів, що склеєні в спільній вершині. Доведено результат типу Амбарцумяна, тобто якщо певна частина спектра така ж сама, як і у випадку нульового потенціалу, то потенціал повинен бути нульовим.

1. Introduction. The addressed problem originates from a work of Ambarzumian [2] on reconstruction of a differential operator from its eigenvalues. Another source of the problem is the so-called *quantum graphs*, i.e., differential operators on graphs [3, 21, 26]. From the classical theory of Sturm – Liouville equations we refer to [7, 17], for special Ambarzumian type inverse problems see [10, 18]. Previous results for graphs are [6, 8, 12, 20, 22, 25]. Both in forward and in inverse problems on graphs a usual ingredient is the calculation of *spectral determinants* (or alternatively *functional determinants* or *characteristic functions*) [1, 5, 9, 11, 13–16, 19, 23, 24, 27]. For a more detailed discussion of these results see the introduction in [20].

2. Results and discussion. Let $r \ge 2$ and consider r cycle graphs C_1, C_2, \ldots, C_r with odd cycle lengths n_1, n_2, \ldots, n_r $(n_j = 1$ is also possible). Let the vertices of C_j be $v_{j0}, \ldots, v_{jn_j} = v_{j0}$, and let us form the graph G as the union of C_j 's, identifying the vertex v_{j0} for all j. We shall say that G is a graph consisting of $r \ge 2$ odd cycles glued together at a common vertex. The edge of G between v_{jk-1} and v_{jk} is sometimes denoted by e_{jk} ; however, when the particular location of the edges are not important, we shall refer to them as $e_1, e_2, \ldots, e_{|E|}$.

Choosing an arbitrary orientation, we parametrize each edge with $x \in [0, 1]$, and consider a Schrödinger operator with potential $q_j(x) \in L^1(0, 1)$ on the edge e_j and with Neumann (or Kirchhoff) boundary conditions (sometimes called standard matching conditions), i.e., solutions are required to be continuous at the vertices and, in the local coordinate pointing outward, the sum of derivatives is zero. More formally, consider the eigenvalue problem

$$-y'' + q_j(x)y = \lambda y \tag{2.1}$$

on e_j for all j with the conditions

$$y_j(\kappa_j) = y_k(\kappa_k) \tag{2.2}$$

if e_j and e_k are incident edges attached to a vertex v where $\kappa = 0$ for outgoing edges, $\kappa = 1$ for incoming edges (and can be both 0 or 1 for loops); and in every vertex v

¹ e-mail: mkiss@math.bme.hu.

² This work was supported by the Hungarian NKFIH (Grant SNN-125119).

M. KISS

$$\sum_{e_j \text{leaves}v} y'_j(0) = \sum_{e_j \text{enters}v} y'_j(1)$$
(2.3)

(loops are counted on both sides).

Theorem 2.1. Consider the eigenvalue problem (2.1)-(2.3). Let G be a graph consisting of $r \ge 2$ odd cycles glued together at a common vertex. If $\lambda = 0$ is the smallest eigenvalue and for infinitely many $k \in \mathbb{Z}^+$ there are r - 1 eigenvalues (counting multiplicities) such that $\lambda = (2k+1)^2 \pi^2 + o(1)$, then q = 0 a.e. on G.

If the lengths of the odd cycles are all 1, i.e., the cycles are all loops, then the statement reduces to that of Theorem 2.1 in [28], which states the following:

Suppose G is a flower-like graph, i.e., a single vertex attached r loops of length 1. For $k = 1, 2, ..., let m_k$ be a sequence of integers with $\lim m_k = +\infty$. If eigenvalues are nonnegative, $\lambda_k = (2m_k + 1)^2 \pi^2$ are eigenvalues with multiplicities (r - 1), where m_k is a strictly ascending infinite sequence of positive integers, then $q_j(x) = 0$ a.e. on [0, 1], for each j = 1, 2, ..., r. We have to require $r \ge 2$ for the consequence to hold.

3. Calculation of the spectral determinant. Denote by $c_j(x, \lambda)$ the solution of (2.1) which satisfies the conditions $c_j(0, \lambda) - 1 = c'_j(0, \lambda) = 0$ and by $s_j(x, \lambda)$ the solution of (2.1) which satisfies the conditions $s_j(0, \lambda) = s'_j(0, \lambda) - 1 = 0$. Each $y_j(x, \lambda)$ may be written as a linear combination

$$y_j(x,\lambda) = A_j(\lambda)c_j(x,\lambda) + \sqrt{\lambda}B_j(\lambda)s_j(x,\lambda).$$

Then $y_j(0,\lambda) = A_j(\lambda)$ is the same on each outgoing edge; hence, as in [20], we index the functions $A(\lambda)$ by vertices, and then

$$y_j(x,\lambda) = A_v(\lambda)c_j(x,\lambda) + \sqrt{\lambda}B_j(\lambda)s_j(x,\lambda),$$

if e_j starts from v. If the eigenfunctions are normalized, i.e., $\sum_j ||y_j(x,\lambda)||_2^2 = 1$, then $A_v(\lambda) = B_j(\lambda) = O(1)$ [8, 20, 28]. The coefficients A_v and B_j form a (|V| + |E|)-dimensional vector, which satisfies |V| Kirchhoff conditions at the vertices and |E| continuity conditions at the incoming ends of edges, namely, for all $v \in V(G)$,

$$\sum_{e_j: \dots \to v} \underbrace{\frac{1}{\sqrt{\lambda}} A_{v_j}(\lambda) c'_j(1,\lambda) + B_j(\lambda) s'_j(1,\lambda)}_{\frac{1}{\sqrt{\lambda}} y'_j(1,\lambda)} - \sum_{e_j: v \to \dots} \underbrace{B_j(\lambda)}_{\frac{1}{\sqrt{\lambda}} y'_j(0,\lambda)} = 0,$$

where in the first sum v_j denotes the starting point of e_j ; and, for all $e_j \in E(G)$,

$$A_u(\lambda)c_j(1,\lambda) + \sqrt{\lambda}B_j(\lambda)s_j(1,\lambda) - A_v(\lambda) = 0,$$

if e_i points from u to v (see equations (2.3) and (2.4) in [20]).

The matrix M of this homogeneous linear system of equations has a special structure. For the convenience of the reader we repeat its description from [20]. Namely, $M = \begin{bmatrix} A & B \\ C & D \end{bmatrix}$, where

A is a |V| by |V| matrix, $a_{vu} = \frac{1}{\sqrt{\lambda}} \sum c'_j(1,\lambda)$, the sum is taken on edges pointing from u to v; in the zero-potential case A is $-\sin\sqrt{\lambda}$ times the (transpose of the) directed adjacency matrix of G;

B and C are like incidence matrices;

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12

1680

$$b_{vj} = \begin{cases} s'_j(1,\lambda), & \text{if } e_j \text{ ends in } v, \\ -1 & \text{if } e_j \text{ starts from } v, \\ s'_j(1,\lambda) - 1, & \text{if } e_j \text{ is a loop in } v, \\ 0, & \text{otherwise} \end{cases}$$

and

$$c_{jv} = \begin{cases} -1, & \text{if } e_j \text{ ends in } v, \\ c_j(1,\lambda), & \text{if } e_j \text{ starts from } v, \\ -1 + c_j(1,\lambda), & \text{if } e_j \text{ is a loop in } v, \\ 0, & \text{otherwise}; \end{cases}$$

D is an |E| by |E| diagonal matrix, $d_{jj} = \sqrt{\lambda} s_j(1, \lambda)$.

The determinant of the matrix M is the so-called spectral determinant of the problem (2.1)–(2.3). *Example.* Consider a flower-like graph, i.e., a single vertex with r loops. Then

$$M = M_1 = \begin{bmatrix} \frac{1}{\sqrt{\lambda}} \sum_{k=1}^r c'_k(1,\lambda) & s'_1(1,\lambda) - 1 & \dots & s'_r(1,\lambda) - 1 \\ \hline -1 + c_1(1,\lambda) & \sqrt{\lambda}s_1(1,\lambda) & \dots & 0 \\ \vdots & \vdots & \ddots & \vdots \\ -1 + c_r(1,\lambda) & 0 & \dots & \sqrt{\lambda}s_r(1,\lambda) \end{bmatrix}$$

with determinant

$$\det M_1 = \lambda^{\frac{r-1}{2}} \left(\sum_{k=1}^r c'_k(1,\lambda) \prod_{j=1}^r s_j(1,\lambda) - \sum_{k=1}^r (s'_k(1,\lambda) - 1)(-1 + c_k(1,\lambda)) \prod_{j \neq k} s_j(1,\lambda) \right)$$

corresponding to formula (2.9) in [28].

For $\lambda = (2k+1)^2 \pi^2 + d + o(1)$ the elements of M have the following asymptotics, independent of k (see [8], equation (2.3), or [22], Lemma 3.1):

$$\frac{1}{\sqrt{\lambda}}c'_j(1,\lambda) = \frac{1}{2\sqrt{\lambda}} \left(d - \int_0^1 q_j \right) + o\left(\frac{1}{\sqrt{\lambda}}\right), \tag{3.1}$$

$$s'_{j}(1,\lambda) = -1 + o\left(\frac{1}{\sqrt{\lambda}}\right),\tag{3.2}$$

$$c_j(1,\lambda) = -1 + o\left(\frac{1}{\sqrt{\lambda}}\right),\tag{3.3}$$

$$\sqrt{\lambda}s_j(1,\lambda) = \frac{1}{2\sqrt{\lambda}} \left(\int_0^1 q_j - d \right) + o\left(\frac{1}{\sqrt{\lambda}}\right).$$
(3.4)

Remark. Using these asymptotics, we get

$$\det M_1 = -4\sum_{k=1}^r \prod_{j \neq k} \sqrt{\lambda} s_j(1,\lambda) + o(\lambda^{-\frac{r}{2}}).$$

This is a special case of (4.1) and of (4.2) below.

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12

1681

4. Proofs of main result.

Lemma 4.1. The total multiplicities of the eigenvalues $\lambda = (2k+1)^2 \pi^2 + O(1)$ are exactly |E| - |V|.

Proof. If q = 0, det M is a polynomial of $\cos \sqrt{\lambda}$ and $\sin \sqrt{\lambda}$, hence its zeros are 2π -periodic in $\sqrt{\lambda}$. $\lambda = (2k+1)^2 \pi^2 + O(1)$ with periodicity implies $\sqrt{\lambda} = (2k+1)\pi$ exactly. For $\lambda = (2k+1)^2 \pi^2$ A and D are zero matrices, thus the rank of M is 2|V|, and its nullspace is exactly (|E| - |V|)-dimensional. Let us write for a moment $\lambda = \lambda(q)$ to denote the dependency of the eigenvalues on the potential. If we arrange the eigenvalues in an increasing sequence $\lambda_1, \ldots, \lambda_n, \ldots$, then there is a constant c depending only on the graph and the L^1 -norm of the potential and not depending on n, the particular index of the eigenvalue, such that $|\lambda_n(q) - \lambda_n(0)| \le c$ for all n [20]. Hence, the total multiplicity of eigenvalues $\lambda = (2k+1)^2 \pi^2 + O(1)$ is the same for all $q \in L^1$.

Lemma 4.2. The determinant of M for $\lambda = (2k+1)^2 \pi^2 + O(1)$ is $O(\lambda^{-\frac{1}{2}(|E|-|V|)})$.

Proof. Each term in the Leibniz formula for the determinant must contain at least (|E| - |V|) factors from A and D having a magnitude of $O\left(\frac{1}{\sqrt{\lambda}}\right)$.

Lemma 4.3. Assume that a graph has the same number of edges as vertices. Then the determinant of its (unoriented) incidence matrix is zero, except if there is no even cycles in the graph and every component contains exactly one (odd) cycle. In that case the determinant is $\pm 2^{\kappa}$ where κ denotes the number of components.

Proof. If the graph contains an even cycle, then the corresponding rows are dependent. If a component contains no cycles, then the corresponding columns are dependent. The number of cycles (including loops) is equal to the number of components, hence, if a component contains more than one cycles or loops, then there must be another component without cycles. In all of these cases the determinant of the incidence matrix is zero. Otherwise it is enough to prove the statement for connected graphs, as the incidence matrix is a direct sum of that of the components. It is true for odd cycles as well as for a single vertex with a loop. If the graph is not a single cycle or a single loop, there is at least one vertex with only one incident edge. Removing this vertex (and its edge) from the graph does not change the absolute value of the determinant of the incidence matrix. This can be repeated until we reach a single cycle or a single loop.

Lemma 4.4. If $\lambda = (2k+1)^2 \pi^2 + O(1)$, then the determinant of a $|V| \times |V|$ submatrix of C (and of B) is $\pm 2^{\kappa} + o\left(\frac{1}{\sqrt{\lambda}}\right)$ if the indices of the rows in C (the columns in B) correspond to the edges of a subgraph which has no even cycles and every component contains exactly one (odd) cycle. Otherwise the determinant is $o\left(\frac{1}{\sqrt{\lambda}}\right)$.

Proof. Leaving out the $o\left(\frac{1}{\sqrt{\lambda}}\right)$ terms from the submatrix we make only $o\left(\frac{1}{\sqrt{\lambda}}\right)$ error in its determinant. What we get is the negative of an incidence matrix of a subgraph with |V| vertices and |V| edges. Then the statement follows from the previous lemma.

Theorem 4.1. If $\lambda = (2k+1)^2 \pi^2 + O(1)$, then

$$\det M = (-1)^{|V|} \sum_{\tau} 4^{\kappa(\tau)} \prod_{e_j \notin \tau} \sqrt{\lambda} s_j(1,\lambda) + O(\lambda^{-\frac{|E| - |V| + 1}{2}}), \tag{4.1}$$

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12

where the sum is taken for such subgraphs τ of G which have |V| vertices and their incidence matrix is nonsingular (i.e., τ has no even cycles, has κ components, each of which contains exactly one (odd) cycle).

Proof. The main terms in the Leibniz formula for the determinant are those which contain exactly (|E| - |V|) elements from D. The product of a fixed set of (|E| - |V|) elements in D is weighted by the determinant of the respective minor, with all other elements of D substituted by zero. The remaining rows in C and columns in B look like an unordered incidence matrix of the graph τ spanned by the remaining |V| edges. Then the determinant of the minor is $(-1)^{|V|}$ times the square of the determinant of the incidence matrix of τ .

Corollary 4.1. If $\lambda = (2k+1)^2 \pi^2 + O(1)$ and the graph G consists of r odd cycles of length n_1, \ldots, n_r , glued together at a common vertex, then

$$\det M = -4\sum_{i=1}^{r} \prod_{j \neq i} \sum_{l=1}^{n_j} \sqrt{\lambda} s(1, \lambda, q_{jl}) + O(\lambda^{-\frac{r}{2}}),$$
(4.2)

where q_{jl} is the potential on the *l*th edge of the *j*th cycle.

Proof. The incidence matrix of a subgraph of G is nonsingular if and only if we leave out one edge from every but one cycle. Note also that |V| is odd and |E| - |V| = r - 1. Then the statement follows from (4.1).

Substituting the asymptotics (3.1)-(3.4), we get the following corollary.

Corollary 4.2. If $\lambda = (2k+1)^2 \pi^2 + d + o(1)$ and the graph G consists of r odd cycles of length n_1, \ldots, n_r , glued together at a common vertex, then

$$\det M = -4\left(\frac{-1}{2\sqrt{\lambda}}\right)^{r-1} p(d) + o(\lambda^{-\frac{r-1}{2}}),$$

where

$$p(d) = \sum_{i=1}^{r} \prod_{j \neq i} \left(n_j d - \sum_{l=1}^{n_j} \int_{0}^{1} q_{jl} \right).$$
(4.3)

Lemma 4.5. Under the assumptions of Theorem 2.1, $p(d) = \sum_{i=1}^{r} \prod_{j \neq i} n_j d^{r-1}$. **Proof.** λ is an eigenvalue of the eigenvalue problem (2.1)–(2.3) if and only if det $M(\lambda) = 0$.

Proof. λ is an eigenvalue of the eigenvalue problem (2.1)-(2.3) if and only if det $M(\lambda) = 0$. Let the distinct roots of p(d) be d_1, \ldots, d_l . By the previous corollary for $\lambda = (2k+1)^2 \pi^2 + O(1)$ the distinct roots of det $M(\lambda)$ are exactly of the form $\lambda = (2k+1)^2 \pi^2 + d_j + o(1), 1 \le j \le l$. By Lemma 4.1 the total multiplicity of these eigenvalues is |E| - |V|. In Theorem 2.1 we assumed that there are the same number of eigenvalues such that $\lambda = (2k+1)^2 \pi^2 + o(1)$, hence $d_j = 0$ for all j and then p(d) is a constant multiple of d^{r-1} . The principal coefficient is given by (4.3).

Proof of Theorem 2.1. Let us introduce $Q_j = \sum_{l=1}^{n_j} \int_0^1 q_{jl}$.

For a fixed m substituting $d = \frac{Q_m}{n_m}$ to (4.3), we get

$$\sum_{i=1}^{r} \frac{1}{n_i} \left(\frac{Q_m}{n_m}\right)^{r-1} = \frac{1}{\prod_{j=1}^{r} n_j} p\left(\frac{Q_m}{n_m}\right) = \frac{1}{n_m} \prod_{j \neq m} \left(\frac{Q_m}{n_m} - \frac{Q_j}{n_j}\right).$$

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12

Introducing $h_j = \frac{Q_j}{n_j}$, we have

$$\sum_{i=1}^{r} \frac{1}{n_i} h_m^{r-1} = \frac{1}{n_m} \prod_{j \neq m} (h_m - h_j), \quad m = 1, 2, \dots, r.$$

We can assume $h_1 \ge h_2 \ge \ldots \ge h_r$. Then, for m = 2, the left-hand side is nonpositive, hence $h_2 \le 0$. Similarly, $h_{r-1} \ge 0$. Hence, for m = 1,

$$\frac{1}{n_1}h_1^{r-2}(h_1 - h_r) = \sum_{i=1}^r \frac{1}{n_i}h_1^{r-1}.$$

If $h_r = 0$, then $h_1 = 0$ as n_j 's are positive. Similarly, $h_1 = 0$ implies $h_r = 0$. If neither of them is zero, then $\frac{1}{n_1}(h_1 - h_r) = \sum_{i=1}^r \frac{1}{n_i}h_1$ and $\frac{1}{n_r}(h_r - h_1) = \sum_{i=1}^r \frac{1}{n_i}h_r$. Subtracting, we get

$$\sum_{i=2}^{r-1} \frac{1}{n_i} (h_1 - h_r) = 0.$$

As n_j 's are positive, if r > 2 then $h_1 = h_2 = \ldots = h_r = 0$, while for r = 2, $h_1n_1 + h_2n_2 = 0$. In both cases,

$$\sum_{j} Q_{j} = \sum_{e_{jl} \in G} \int_{0}^{1} q_{jl} = \int_{G} q = 0.$$

Let us denote the operator of the eigenvalue problem (2.1)–(2.3) by L. $\frac{\langle \varphi, L\varphi \rangle}{\langle \varphi, \varphi \rangle} \ge \lambda_0 = 0$ and equality holds if and only if φ is an eigenfunction of L. It follows that the constant 1 must be an eigenfunction corresponding to the eigenvalue 0. Substituting this to (2.1) gives q(x) = 0.

Theorem 2.1 is proved.

5. Glossary. A walk W in a graph is an alternating sequence of vertices and edges, say $X_0, e_1, \ldots, e_l, X_l$ where $e_i = X_{i-1}X_i$, 0 < i < l. The length of W is l. This walk W is called a *trail* if all its edges are distinct. A *path* is a walk with distinct vertices. A trail whose end vertices coincide (a closed trail) is called a *circuit*. To be precise, a circuit is a closed trail without distinguished endvertices and direction, so that, for example, two triangles sharing a single vertex give rise to precisely two circuits with six edges. If a walk $W = X_0, e_1, \ldots, e_l, X_l$ is such that $l \ge 3, X_0 = X_l$, and the vertices $X_i, 0 < i < l$, are distinct from each other and X_0 , then W is said to be a *cycle* [4, p. 5].

The *incidence matrix* of a graph has a row for each vertex and a column for each edge, and is defined as

$$R = (r_{ij}), \qquad r_{ij} = \begin{cases} 0, & \text{if } e_j & \text{is not incident to } v_i, \\ 1, & \text{if } e_j & \text{is not a loop and incident to } v_i, \\ 2, & \text{if } e_j & \text{is a loop at } v_i. \end{cases}$$

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12

1684

References

- 1. E. Akkermans, A. Comtet, J. Desbois, G. Montambaux, Ch. Texier, *Spectral determinant on quantum graphs*, Annal. Phys., **284**, № 1, 10–51 (2000).
- 2. V. Ambarzumian, Über eine Frage der Eigenwerttheorie, Z. Phys., 53, 690-695 (1929).
- G. Berkolaiko, P. Kuchment, *Introduction to quantum graphs*, Math. Surveys and Monogr., Amer. Math. Soc. (2013).
 B. Bollobás, *Modern graph theory*, Springer, New York (1998).
- 5. J. Bolte, S. Egger, R. Rueckriemen, *Heat-kernel and resolvent asymptotics for Schrödinger operators on metric graphs*, Appl. Math. Res. Express, **2015**, № 1, 129–165 (2014).
- 6. J. Boman, P. Kurasov, R. Suhr, Schrödinger operators on graphs and geometry II. Spectral estimates for L₁-potentials and an Ambartsumian theorem, Integral Equat. and Operator Theory, **90**, № 3 (2018).
- 7. G. Borg, Eine Umkehrung der Sturm-Liouvilleschen Eigenwertaufgabe, Acta Math., 78, 1–96 (1946).
- R. Carlson, V. Pivovarchik, Ambarzumian's theorem for trees, Electron. J. Different. Equat., 2007, № 142, 1-9 (2007).
- 9. R. Carlson, V. Pivovarchik, *Spectral asymptotics for quantum graphs with equal edge lengths*, J. Phys. A: Math. and Theor., **41**, № 14, Article 145202 (2008).
- Y. H. Cheng, Tui-En Wang, Chun-Jen Wu, A note on eigenvalue asymptotics for Hill's equation, Appl. Math. Lett.,
 23, № 9, 1013-1015 (2010).
- S. Currie, B. A. Watson, *Eigenvalue asymptotics for differential operators on graphs*, J. Comput. and Appl. Math., 182, № 1, 13–31 (2005).
- 12. E. B. Davies, An inverse spectral theorem, J. Operator Theory, 69, № 1, 195-208 (2013).
- 13. Yves Colin De Verdière, *Semi-classical measures on quantum graphs and the Gauss map of the determinant manifold*, Ann. Henri Poincaré, **16**, 347 364 (2015).
- 14. J. Desbois, Spectral determinant of Schrödinger operators on graphs, J. Phys. A: Math. and Gen., 33, № 7 (2000).
- 15. L. Friedlander, Determinant of the Schrödinger operator on a metric graph, Contemp. Math., 415, 151-160 (2006).
- 16. J. M. Harrison, K. Kirsten, C. Texier, Spectral determinants and zeta functions of Schrödinger operators on metric graphs, J. Phys. A: Math. and Theor., 45, № 12, Article 125206 (2012).
- 17. M. Horváth, Inverse spectral problems and closed exponential systems, Ann. Math., 162, № 2, 885–918 (2005).
- 18. M. Horváth, On the stability in Ambarzumian theorems, Inverse Problems, 31, № 2, Article 025008 (2015).
- 19. I. Kac, V. Pivovarchik, On multiplicity of a quantum graph spectrum, J. Phys. A: Math. and Theor., 44, № 10, Article 105301 (2011).
- 20. M. Kiss, *Spectral determinants and an Ambarzumian type theorem on graphs*, Integral Equat. and Operator Theory, **92**, 1–11 (2020).
- 21. P. Kuchment, *Quantum graphs: an introduction and a brief survey*, Analysis on Graphs and its Applications, 291–314 (2008).
- 22. Chun-Kong Law, Eiji Yanagida, *A solution to an Ambarzumyan problem on trees*, Kodai Math. J., **35**, № 2, 358–373 (2012).
- 23. M. Möller, V. Pivovarchik, Spectral theory of operator pencils, Hermite Biehler functions, and their applications, Oper. Theory: Adv. and Appl., Springer (2015).
- 24. K. Pankrashkin, Spectra of Schrödinger operators on equilateral quantum graphs, Lett. Math. Phys., 77, № 2, 139–154 (2006).
- 25. V. N. Pivovarchik, Ambarzumian's theorem for a Sturm Liouville boundary value problem on a star-shaped graph, Funct. Anal. and Appl., **39**, № 2, 148–151 (2005).
- 26. Yu. V. Pokornyi, A. V. Borovskikh, *Differential equations on networks (geometric graphs)*, J. Math. Sci., **119**, № 6, 691-718 (2004).
- 27. Ch. Texier, ζ-Regularized spectral determinants on metric graphs, J. Phys. A: Math. and Theor., 43, № 42, Article 425203 (2010).
- 28. Chuan-Fu Yang, Xiao-Chuan Xu, Ambarzumyan-type theorems on graphs with loops and double edges, J. Math. Anal. and Appl., 444, № 2, 1348-1358 (2016).

Received 10.05.21

ISSN 1027-3190. Укр. мат. журн., 2022, т. 74, № 12